

String theory, sheaf cohomology and the *homalg* package

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Outline

- 1 Brief introduction to string theory

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- 2 Sheaf cohomology in string theory and the homalg package

Section 1

Brief introduction to string theory

Getting started with string theory

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Why string theory?

- Current understanding of physics:

General relativity + Quantum field theory

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⇒ String theory is a good candidate theory

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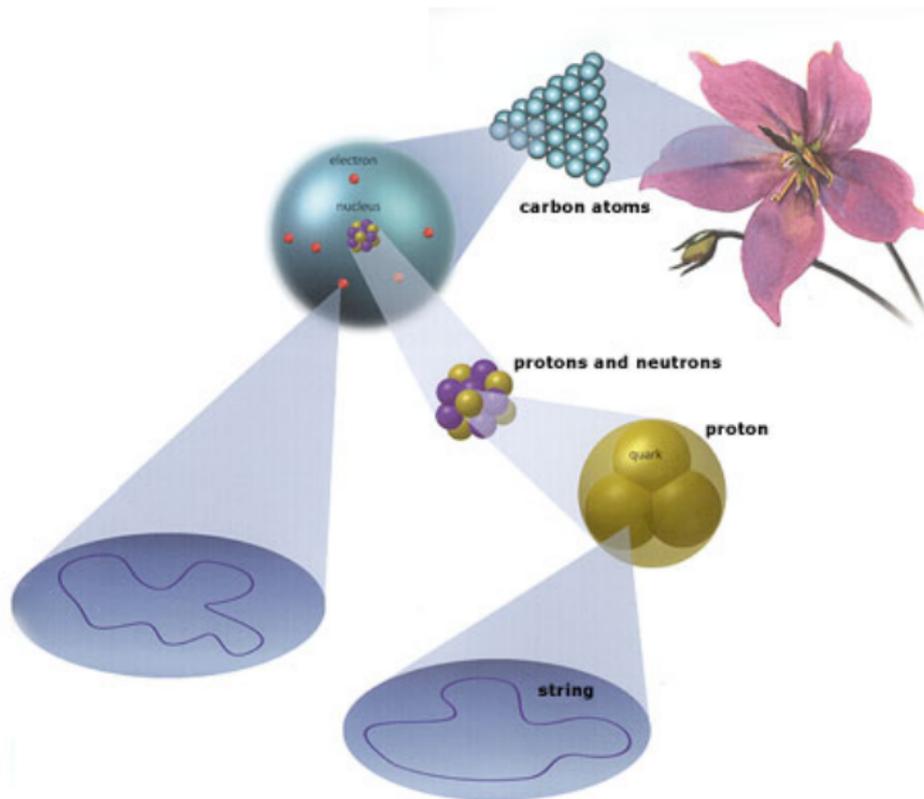
⇒ String theory is a good candidate theory

Starting point of string theory:

Replace point particles by 1-dimensional objects.

Replacing point particles by strings

from 'a Layman's guide to string theory'



Minor change with big consequences

Spacetime

- We are living in a 10d Riemannian manifold \mathcal{S}
- Compactification means $\mathcal{S} = \mathbb{R}^{1,3} \times_w \mathcal{M}_6$ and \mathcal{M}_6 **compact**

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- Vibrations of string $\hat{=}$ elementary particles.
- Every (closed) string gives rise to $|\mathbb{N}|$ elementary particles $|n\rangle$ with masses

$$M^2(|n\rangle) = \frac{4}{\alpha'} (n - 1)$$

where $\alpha' \ll 1$ is the *Regge slope*.

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- ⇒ Eventually focus on the **massless** particles

Further consequence - D-branes appear in string theory

Rough picture

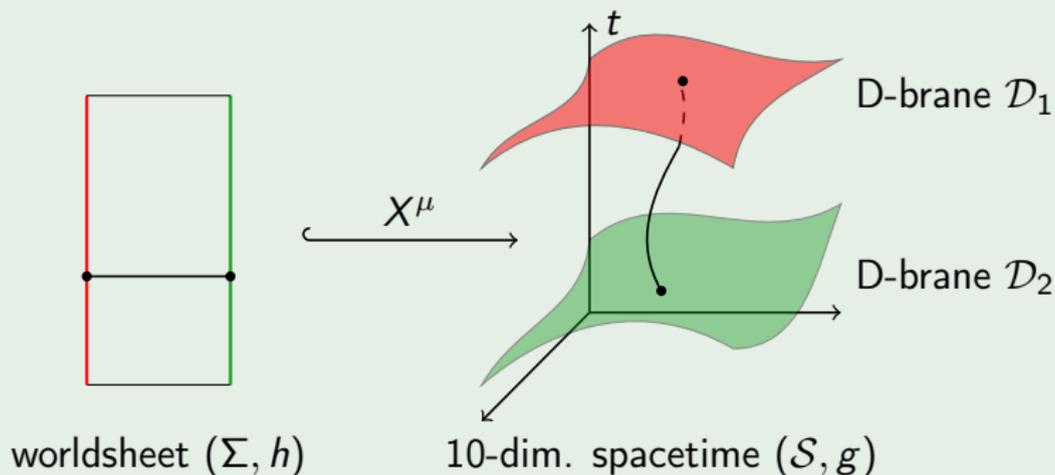
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D-branes are places where open strings end

Picture of open strings and D-branes



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- Each solution corresponds to the physics of an entire universe

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- A number of solutions come quite close to our universe.
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⇒ An enormous search!

The search in the string landscape



Section 2

Sheaf cohomology in string theory and the homalg package

Special solutions to string theory: the B-model hep-th/9112056

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The B-model ...

is the region of the string landscape described by hep-th/0403166

- $\mathcal{S} = \mathbb{R}^{1,3} \times \mathcal{M}_6$ where \mathcal{M}_6 is a CY 3-fold
- use $\mathcal{N} = (2, 2)$ supersymmetric theory on \mathcal{M}_6
- string coupling $g_s \rightarrow 0$
- focus on topological sector
- ...

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Dictionary in B-model hep-th/0208104

Physics	Mathematics
D-brane \mathcal{D} on \mathcal{M}_6	$\mathcal{D} \in D^b(\mathcal{Coh}(\mathcal{M}_6))$
Massless string excitation s between D-branes $\mathcal{F}_\bullet, \mathcal{G}_\bullet$	$s \in \text{Ext}^q(\mathcal{F}_\bullet, \mathcal{G}_\bullet)$

Special solutions to string theory: F-theory GUT model

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- 1 $X_\Sigma = \mathbb{CP}^3 \times \mathbb{CP}^1$,
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- 3 $S_{\text{GUT}} \subset B_3$ with $\dim S_{\text{GUT}} = 2$

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 $V(\langle x_0, x_0^3 + x_1^3 + x_2^3 + x_3^3 \rangle)$

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Step 2a: Elliptic fibration over B_3 - General Story

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- For $i \in \{1, 2, 3, 4, 6\}$ pick $a_i \in H^0(X_\Sigma, \mathcal{O}_{X_\Sigma}(i\bar{K}_{B_3}))$.

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- For $[x, y, z] \in \mathbb{CP}_{2,3,1}$, $p \in B_3$ set

$$\mathcal{C}_p := \{[x, y, z] \in \mathbb{CP}_{2,3,1}, Q(x, y, z, p) = 0\}$$

$$Q(x, y, z, p) := x^2 - y^2 + xyza_1(p) + x^2z^2a_2(p) + yz^3a_3(p) \\ + xz^4a_4(p) + z^6a_6(p)$$

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- $Y_4 := \bigcup_{p \in B_3} C_p$ with $\pi: Y_4 \rightarrow B_3$ s.t. $\pi^{-1}(p) = C_p$

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Step 2b: Elliptic fibration over B_3 - Our example

- $X_\Sigma = \mathbb{CP}^3 \times \mathbb{CP}^1$, $S = \mathbb{C}[x_0, \dots, x_3, y_0, y_1]$
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Step 3: Build A $SU(5) \times U(1)_x$ Model Over S_{GUT}

Require that $a_2 = a_{2,1} w$, $a_3 = a_{3,2} w^2$, $a_4 = a_{4,3} w^3$, $a_6 \equiv 0$ s.t.

- $a_{j,k} \in H^0[X_\Sigma, \mathcal{O}_{X_\Sigma}(j \cdot \bar{K}_{B_3} - k \cdot D_{\text{GUT}})]$
- $a_1, a_{j,k}$ not divisible by w in S

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- $C_{10} := \{p \in X_\Sigma, s_{B_3} = w = a_1 = 0\}$
- $C_{\bar{5}_m} := \{p \in X_\Sigma, s_{B_3} = w = a_{3,2} = 0\}$
- $C_{5_H} := \{p \in X_\Sigma, s_{B_3} = w = a_{3,2} \cdot a_{2,1} - a_{4,3} \cdot a_1 = 0\}$

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Step 5: Add fluxes and compute (massless) particle spectrum

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- ⇒ The actual search is then a scan over admissible fluxes $(\mathcal{F}, \mathcal{H})$.

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Further improvements needed:

As much performance as possible - there is a big search ahead.

Thank you for your attention! Questions?

